

USANDO

$$H^{\mu\nu} = \frac{i}{2} (x^\mu \partial^\nu - x^\nu \partial^\mu)$$

VALUTIAMO

$$[H^{\mu\nu}, H^{\rho\sigma}]$$

PER NON

PERDERE

PEZZI

SERVIAMOCI

DI UNA

FUNZIONE

$$f = f(x^\mu);$$

$$\begin{aligned} H^{\mu\nu} H^{\rho\sigma} f &= -\frac{1}{2} (x^\mu \partial^\nu - x^\nu \partial^\mu) (x^\rho \partial^\sigma - x^\sigma \partial^\rho) f \\ &= -\frac{1}{2} [x^\mu \partial^\nu (x^\rho \partial^\sigma) - x^\mu \partial^\nu (x^\sigma \partial^\rho) - x^\nu \partial^\mu (x^\rho \partial^\sigma) + x^\nu \partial^\mu (x^\sigma \partial^\rho)] f(x) \\ &= -\frac{1}{2} [x^\mu (\delta^{\nu\rho} \partial^\sigma + x^\rho \partial^\nu \partial^\sigma) - x^\mu (\delta^{\nu\sigma} \partial^\rho + x^\sigma \partial^\nu \partial^\rho) - \\ &\quad - x^\nu (\delta^{\mu\rho} \partial^\sigma + x^\rho \partial^\mu \partial^\sigma) + x^\nu (\delta^{\mu\sigma} \partial^\rho + x^\sigma \partial^\mu \partial^\rho)] \end{aligned}$$

$$\begin{aligned} \partial^\nu &= g^{\nu\alpha} \partial_\alpha \\ \partial^\nu x^\rho &= g^{\nu\alpha} \partial_\alpha x^\rho \\ &= g^{\nu\alpha} \delta_\alpha^\rho \\ &= g^{\nu\rho} \end{aligned}$$

IN QUESTO MODO:

$$[H^{\mu\nu}, H^{\rho\sigma}] f = [H^{\mu\nu} H^{\rho\sigma} - H^{\rho\sigma} H^{\mu\nu}] f =$$

$$\begin{aligned} &= -\frac{1}{2} [x^\mu \delta^{\nu\rho} \partial^\sigma + x^\mu x^\rho \partial^\nu \partial^\sigma - x^\mu \delta^{\nu\sigma} \partial^\rho - x^\mu x^\sigma \partial^\nu \partial^\rho - x^\nu \delta^{\mu\rho} \partial^\sigma - x^\nu x^\rho \partial^\mu \partial^\sigma + \\ &\quad + x^\nu \delta^{\mu\sigma} \partial^\rho + x^\nu x^\sigma \partial^\mu \partial^\rho] f + [x^\rho \delta^{\sigma\mu} \partial^\nu + x^\rho x^\mu \partial^\sigma \partial^\nu - x^\rho \delta^{\sigma\nu} \partial^\mu - x^\rho x^\nu \partial^\sigma \partial^\mu - \\ &\quad - x^\sigma \delta^{\rho\mu} \partial^\nu - x^\sigma x^\mu \partial^\rho \partial^\nu + x^\sigma \delta^{\rho\nu} \partial^\mu + x^\sigma x^\nu \partial^\rho \partial^\mu] f \end{aligned}$$

VISTA LA  
SIMMETRIA  
DI  $g^{\mu\nu}$ TERMINI  
RACCOLGERE:

SOTTOLINEATI A ZIG-ZAG SI CANCELLANO, MENTRE GLI ALTRI LI POSSIAMO

$$[H^{\mu\nu}, H^{\rho\sigma}] = -\frac{1}{2} [g^{\nu\rho} (x^\mu \partial^\sigma - x^\sigma \partial^\mu) + g^{\nu\sigma} (x^\rho \partial^\mu - x^\mu \partial^\rho)] + \dots$$

(1)

$$\dots + g^{\mu\rho}(\underline{x^\sigma \partial^\nu - x^\nu \partial^\sigma}) + g^{\mu\sigma}(\underline{x^\nu \partial^\rho - x^\rho \partial^\nu})]$$

POSSIAMO RICONOSCERE :  $-\dot{\lambda} H^{\mu\nu} = (x^\mu \partial^\nu - x^\nu \partial^\mu)$

$$[H^{\mu\nu}, H^{\rho\sigma}] = -\dot{\lambda} [g^{\nu\rho} H^{\mu\sigma} + g^{\nu\sigma} H^{\rho\mu} + g^{\mu\rho} H^{\sigma\nu} + g^{\mu\sigma} H^{\nu\rho}]$$

POSSIAMO ANCHE SISTEMARE I SEGNI USANDO LA SIMMETRIA DI  $g^{\mu\nu}$  E L'ANTISIMMETRIA DI  $H^{\mu\nu}$ . LA STRUTTURA DEGLI INDICI DEI TERMI È:

$$[12, 34] = -\dot{\lambda} [(14, 23) + (23, 14) - (13, 24) - (24, 13)]$$

$\swarrow$  INDICI  $g$        $\searrow$  INDICI  $H$

SISTEMANDO LA NOSTRA SCRITTURA PER OTTENERE:

$$[H^{\mu\nu}, H^{\rho\sigma}] = -\dot{\lambda} [(\mu\sigma, \nu\rho) + (\nu\rho, \mu\sigma) - (\mu\rho, \nu\sigma) - (\nu\sigma, \mu\rho)]$$

NOI ABBIAMO

$$= -\dot{\lambda} [g^{\mu\sigma} H^{\nu\rho} + g^{\nu\rho} H^{\mu\sigma} - g^{\mu\rho} H^{\nu\sigma} - g^{\nu\sigma} H^{\mu\rho}]$$

AVENDO

$$N_i = \frac{1}{2}(\mathcal{J}_i + i K_i)$$

$$N_i^\dagger = \frac{1}{2}(\mathcal{J}_i - i K_i)$$

E

FACILE

VERIFICARE

VALGONO!

$$[\mathcal{J}_i, \mathcal{J}_j] = i \varepsilon_{ijk} \mathcal{J}_k$$

$$[K_i, K_j] = -i \varepsilon_{ijk} \mathcal{J}_k$$

$$[\mathcal{J}_i, K_j] = i \varepsilon_{ijk} K_k$$

$$[N_i, N_j^\dagger] = \left[ \frac{1}{2}(\mathcal{J}_i + i K_i), \frac{1}{2}(\mathcal{J}_j - i K_j) \right]$$

$$= \frac{1}{4} \left( [\mathcal{J}_i, \mathcal{J}_j] - i [\mathcal{J}_i, K_j] + i [K_i, \mathcal{J}_j] + [K_i, K_j] \right)$$

$$= \frac{1}{4} \left( i \varepsilon_{ijk} \mathcal{J}_k + i \varepsilon_{ijk} K_k + i \underbrace{(-i \varepsilon_{ijk} K_k)}_{-[\mathcal{J}_j, K_i] = i \varepsilon_{jki} \mathcal{J}_k} - i \varepsilon_{ijk} \mathcal{J}_k \right)$$

$$= \frac{1}{4} \left( i \varepsilon_{ijk} \mathcal{J}_k + \varepsilon_{ijk} K_k - \varepsilon_{ijk} K_k - i \varepsilon_{ijk} \mathcal{J}_k \right)$$

$$= 0$$

$$\begin{aligned} [A, B] &= AB - BA \\ [B, A] &= BA - AB \\ &= -[A, B] \end{aligned}$$

$$[N_i, N_j] = \frac{1}{4} [\mathcal{J}_i + i K_i, \mathcal{J}_j + i K_j] = \frac{1}{4} \left( [\mathcal{J}_i, \mathcal{J}_j] + i [\mathcal{J}_i, K_j] + i [K_i, \mathcal{J}_j] - [K_i, K_j] \right)$$

$$= \frac{1}{4} \left( i \varepsilon_{ijk} \mathcal{J}_k + i (\varepsilon_{ijk} K_k) + i \underbrace{(-i \varepsilon_{jki} K_k)}_{-[\mathcal{J}_j, K_i]} + \varepsilon_{ijk} \mathcal{J}_k \right)$$

$$= \frac{1}{2} i \varepsilon_{ijk} \mathcal{J}_k - \frac{1}{2} \varepsilon_{ijk} K_k$$

$$= i \varepsilon_{ijk} \left( \frac{1}{2} (\mathcal{J}_k + i K_k) \right) = i \varepsilon_{ijk} N_k$$



$$[N_i^+, N_j^+] \cdot \frac{1}{4} [\mathcal{J}_i - iK_i, \mathcal{J}_j - iK_j]$$

$$= \frac{1}{4} \left( [\mathcal{J}_i, \mathcal{J}_j] - i[\mathcal{J}_i, K_j] - i[K_i, \mathcal{J}_j] - [K_i, K_j] \right)$$

$$\underbrace{- [\mathcal{J}_j, K_i]}_{= -i\varepsilon_{jik} K_k = i\varepsilon_{ijk} K_k}$$

$$= \frac{1}{4} \left( i\varepsilon_{ijk} \mathcal{J}_k - i(i\varepsilon_{ijk} K_k) - i(i\varepsilon_{ijk} K_k) + i\varepsilon_{ijk} K_k \right)$$

$$= \frac{1}{2} i\varepsilon_{ijk} \mathcal{J}_k + \frac{1}{2} \varepsilon_{ijk} K_k$$

$$= i\varepsilon_{ijk} \frac{1}{2} (\mathcal{J}_k - iK_k)$$

$$= i\varepsilon_{ijk} N_k^+$$

CONOSCO

$$(M^{\mu\nu})^\mu{}_\nu = -i(g^{\mu\beta}\delta^\nu{}_\nu - g^{\mu\nu}\delta^\beta{}_\nu)$$

POSSIAMO DETERMINARE LE 6 MATRICI  $M^{\mu\nu}$ .

$$(M^{01})^\mu{}_\nu = -i(g^{\mu 0}\delta^1{}_\nu - g^{\mu 1}\delta^0{}_\nu) = -i\left[\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}(0\ 1\ 0\ 0) - \begin{pmatrix} 0 \\ -1 \\ 0 \\ 0 \end{pmatrix}(1\ 0\ 0\ 0)\right] = -i\begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{02})^\mu{}_\nu = -i(g^{\mu 0}\delta^2{}_\nu - g^{\mu 2}\delta^0{}_\nu) = -i\left[\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}(0\ 0\ 1\ 0) - \begin{pmatrix} 0 \\ 0 \\ -1 \\ 0 \end{pmatrix}(1\ 0\ 0\ 0)\right] = -i\begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{03})^\mu{}_\nu = -i(g^{\mu 0}\delta^3{}_\nu - g^{\mu 3}\delta^0{}_\nu) = -i\left[\begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}(0\ 0\ 0\ 1) - \begin{pmatrix} 0 \\ 0 \\ 0 \\ -1 \end{pmatrix}(1\ 0\ 0\ 0)\right] = -i\begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{12})^\mu{}_\nu = -i(g^{\mu 1}\delta^2{}_\nu - g^{\mu 2}\delta^1{}_\nu) = -i\left[\begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}(0\ 0\ 1\ 0) - \begin{pmatrix} 0 \\ 0 \\ -1 \\ 0 \end{pmatrix}(0\ 1\ 0\ 0)\right] = -i\begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{13})^\mu{}_\nu = -i(g^{\mu 1}\delta^3{}_\nu - g^{\mu 3}\delta^1{}_\nu) = -i\left[\begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix}(0\ 0\ 0\ 1) - \begin{pmatrix} 0 \\ 0 \\ 0 \\ -1 \end{pmatrix}(0\ 0\ 1\ 0)\right] = -i\begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & +1 & 0 \end{pmatrix}$$

$$(M^{23})^\mu{}_\nu = -i(g^{\mu 2}\delta^3{}_\nu - g^{\mu 3}\delta^2{}_\nu) = -i\left[\begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}(0\ 0\ 0\ 1) - \begin{pmatrix} 0 \\ 0 \\ 0 \\ -1 \end{pmatrix}(0\ 1\ 0\ 0)\right] = -i\begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

(5)

POSSIAMO VERIFICARE CHE

$$J_i = \frac{1}{2} \varepsilon_{ijk} M^{jk}$$
 (UGUALE A  $J_i = \frac{1}{2} \varepsilon_{ijk} M^{jk}$  POICHÈ ABBASSANDO 2)  
 INDICI SPAZIALI ESCE "(-)(-) = +"

$$K^i = M^{0i}$$

È IMMEDIATO

$$K^1 = K_x = M^{01} = -\dot{\lambda} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$K^2 = K_y = M^{02} = -\dot{\lambda} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

VEDI  
ILIPOPOULOS  
PAG. 94

$$K^3 = K_z = M^{03} = -\dot{\lambda} \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

$$J^1 = J^x = \frac{1}{2} \varepsilon^{ijk} M^{jk} = \frac{1}{2} \left( \varepsilon^{123} M^{23} + \underbrace{\varepsilon^{132} M^{32}}_{(-\varepsilon^{123})(-M^{23})} \right) = \underbrace{\varepsilon^{123}}_{+1} M^{23} = -\dot{\lambda} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & +1 & 0 \end{pmatrix}$$

$$J^2 = J^y = \frac{1}{2} \varepsilon^{2jk} M^{jk} = \frac{1}{2} \left( \underbrace{\varepsilon^{213} M^{13}}_{-\varepsilon^{123}} + \underbrace{\varepsilon^{231} M^{31}}_{\varepsilon^{123}(-M^{13})} \right) = -(M^{13}) = -\dot{\lambda} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & +1 \\ 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix}$$

$$J^3 = J^z = \frac{1}{2} \varepsilon^{3jk} M^{jk} = \frac{1}{2} \left( \varepsilon^{312} M^{12} + \underbrace{\varepsilon^{321} M^{21}}_{(-\varepsilon^{123})(-M^{12})} \right) = M^{12} = -\dot{\lambda} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

(6)



VEDIAMO I MOMENTI ANGOLARI  $J_i$  E I BOOST  $K_i$  COSA DIVENTANO PER TRASFORMAZIONI DI PARITÀ.

VEDIAMO AD ESEMPIO  $J_1$ :

$$\rightarrow \Lambda_P = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

VEDI §3.7.2  
DI SCHWICHTENBERG

$$\begin{aligned} J_1' &= \Lambda_P J_1 \Lambda_P^\dagger = \Lambda_P J_1 \Lambda_P^T \\ &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & \lambda & 0 \\ 0 & 0 & -\lambda & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}^T \\ &= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & \lambda & 0 \\ 0 & 0 & -\lambda & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \\ &= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & \lambda & 0 \\ 0 & 0 & -\lambda & 0 \end{pmatrix} = J_1 \end{aligned}$$

ANALOGHI CONTI PER  $J_2$  E  $J_3$ . PER IL BOOST  $K_1$ :

$$\begin{aligned} K_1' &= \Lambda_P K_1 \Lambda_P^T = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & -\lambda & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}^T \\ &= \begin{pmatrix} 0 & +\lambda & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} = -K_1 \end{aligned}$$

CONTI ANALOGHI PER  $K_2$ ,  $K_3$ .

CON CONTI ANALOGHI SI VEDE PER L'INVERSIONE TEMPORALE:

$$J_i \xrightarrow{T} J'_i = \Lambda_T J_i \Lambda_T^T = J_i$$

$$K_i \xrightarrow{T} K'_i = \Lambda_T K_i \Lambda_T^T = -K_i$$

QUINDI COMPLESSIVAMENTE SCRIVIAMO

$$J_i \xrightarrow{P} J_i$$

$$K_i \xrightarrow{P} -K_i$$

$$J_i \xrightarrow{T} J_i$$

$$K_i \xrightarrow{T} -K_i$$



$$\Lambda_{LR} = e^{+\frac{i}{2}\sigma^i(\omega_i \mp \eta_i)}$$

$$\Lambda_{L,R}^* = e^{-\frac{i}{2}\sigma^i(\omega_i \pm \eta_i)}$$

$$\Lambda_{L,R}^T = e^{+\frac{i}{2}\sigma^{iT}(\omega_i \mp \eta_i)}$$

$$\Lambda_{L,R}^{-1} = e^{-\frac{i}{2}\sigma^i(\omega_i \mp \eta_i)}$$

$$\Lambda_{L,R}^+ = e^{-\frac{i}{2}\sigma^{i+}(\omega_i \pm \eta_i)}$$

$$\begin{cases} (\sigma^2)^* = -\sigma^2 & ; (\sigma^2)^T = -\sigma^2 \\ (\sigma^2)^T = \sigma^2 & ; (\sigma^i)^+ = \sigma^i \end{cases}$$

$$\sigma^2 \sigma^{-1} \sigma^2 = -\sigma^{i*} \Rightarrow -\sigma^i = \sigma^2 \sigma^{i*} \sigma^2$$

$$\textcircled{1} \Lambda_{L,R}^* \xrightarrow{T} \Lambda_{L,R}^+ = e^{-\frac{i}{2}\sigma^i(\omega_i \pm \eta_i)}$$

$$\Rightarrow \boxed{\Lambda_{L,R}^+ = \Lambda_{R,L}^{-1}}$$

$$\sigma^2 \sigma^2$$

$$\textcircled{3} \Lambda_{L,R}^* = \sigma^2 \Lambda_{R,L} \sigma^2$$

$$\textcircled{2} \boxed{\sigma^2 \Lambda_{L,R}^* \sigma^2 = e^{+\frac{i}{2}\sigma^i(\omega_i \pm \eta_i)} = \Lambda_{R,L}}$$

$$\textcircled{4} \textcircled{3} \xrightarrow{+ = * + T} (\Lambda_{L,R}^*)^+ = (\sigma^2 \Lambda_{R,L} \sigma^2)^+ \Rightarrow \Lambda_{L,R}^T = \sigma^2 \Lambda_{R,L}^+ \sigma^2 \xRightarrow{\textcircled{1}} \boxed{\Lambda_{L,R}^T = \sigma^2 \Lambda_{L,R}^{-1} \sigma^2}$$

$$\textcircled{5} \boxed{\Lambda_L^T \sigma^2 \Lambda_L = \Lambda_L^T \sigma^2 \Lambda_L \sigma^2 \sigma^2 = \Lambda_L^T \Lambda_R^* \sigma^2 = \sigma^2 \Lambda_L^{-1} \sigma^2 \Lambda_R^* \sigma^2 = \sigma^2 \Lambda_L^{-1} \Lambda_L = \sigma^2} \quad \textcircled{4} \quad \textcircled{3}$$

9

CONTO  
GUARDO

GENERICO

PER

$\phi(x^\mu)$

$\mathcal{L}(\phi, \partial_\mu \phi)$

$$\Delta S = \int d^4x \Delta \mathcal{L} + \int \delta(d^4x) \mathcal{L}$$

$$\text{so } \delta(d^4x) = \det \delta x^\mu d^4x$$

PER UNA TRASFORMAZIONE  
DI COORDINATE (VEG)  
NOTE)

CALCOLO

$$\begin{aligned} \Delta \mathcal{L} &= \delta \mathcal{L} + \delta x^\mu \partial_\mu \mathcal{L} \\ &= \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta (\partial_\mu \phi) + \delta x^\mu \partial_\mu \mathcal{L} \\ &= \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \partial_\mu (\delta \phi) + \delta x^\mu \partial_\mu \mathcal{L} \\ &= \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi \right) - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi + \delta x^\mu \partial_\mu \mathcal{L} \end{aligned}$$

IL NETTO ON-SHELL  $\Rightarrow$  EQ MOTO  $= 0$

$$\Rightarrow \Delta \mathcal{L} = \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi \right) + \delta x^\mu \partial_\mu \mathcal{L}$$

COSÌ

$$\Delta S = \int d^4x \left[ \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi \right) + \underbrace{\delta x^\mu \partial_\mu \mathcal{L} + \partial_\mu \delta x^\mu \mathcal{L}}_{\partial_\mu (\delta x^\mu \mathcal{L})} \right]$$

$$\Rightarrow \Delta S = \int d^4x \underbrace{\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi + \delta x^\mu \mathcal{L} \right)}_{J^\mu} = 0$$

QUINDI ABBIAMO

$$J^\mu = \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi + \delta x^\mu \mathcal{L}$$

SE DIFFERENZIO IL CASO DI TRASFORMAZIONE DI COORDINATE:

$$x^\mu \rightarrow x^\mu + \delta x^\mu \Rightarrow \Delta \phi = 0 \Rightarrow \delta \phi = -\delta x^\mu \partial_\mu \phi$$

SOSTITUISCO

$$J^\mu = -\frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta x^\nu \partial_\nu \phi + \delta x^\mu \mathcal{L} = \left( -\frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \partial_\nu \phi + g^\mu{}_\nu \mathcal{L} \right) \delta x^\nu = \boxed{-T^\mu{}_\nu \delta x^\nu}$$

TENSORE  
ENERGIA-IMPULSO

VEDIAMO ANCHE IL CASO GENERICO:

$$x^\mu \rightarrow x^\mu + \delta x^\mu, \quad \phi \rightarrow \phi + \alpha \Rightarrow \Delta \phi = \delta \phi + \delta x^\mu \partial_\mu \phi$$

SOSTITUENDO  ~~$\delta \phi$~~   $\delta \phi = \Delta \phi - \delta x^\mu \partial_\mu \phi$

$$J^\mu = -T^\mu{}_\nu \delta x^\nu + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \Delta \phi$$



$$\delta S = S(x') - S(x) = \int d^4x (L'(x) - L(x)) = \int d^4x \delta L$$

$$d^4x' = (1 + \partial_\mu \delta x^\mu) d^4x$$

$$\Delta S = S(x') - S(x) = \int d^4x' L'(x') - \int d^4x L(x) \quad L'(x) = \Delta L + L(x)$$

$$= \int (1 + \partial_\mu \delta x^\mu) (\Delta L + L(x)) d^4x - \int d^4x L(x)$$

$$\approx \int d^4x (\Delta L + \partial_\mu \delta x^\mu L) = \int d^4x (\delta L + \underbrace{\delta x^\mu \partial_\mu L + L \partial_\mu \delta x^\mu}_{\Delta L = \delta L + \delta x^\mu \partial_\mu L})$$

$$\Delta S = \delta S + \int d^4x 2 \left( \frac{\partial L}{\partial \delta x^\mu} \right)$$

(12)

DIMOSTRAZIONE

$$h = \frac{\vec{J} \cdot \vec{P}}{|\vec{P}|}$$

PARTIAMO DAL SAPERE:

$$W^\mu = \frac{1}{2} \varepsilon^{\mu\nu\sigma\tau} P_\nu M_{\sigma\tau}$$

$$\begin{cases} W^0 = \frac{1}{2} \varepsilon^{0ijk} P_i M_{jk} = P_i J^i = \vec{P} \cdot \vec{J} \\ \vec{W} = P_0 \vec{J} - \vec{P} \times \vec{K} \end{cases}$$

PRENDIAMO  $m=0 \Rightarrow P^\mu = (0, \vec{P})$

$$\Rightarrow \begin{cases} P_\mu P^\mu = 0 \\ W_\mu W^\mu = 0 \\ W_\mu P^\mu = 0 \end{cases} \Rightarrow \boxed{W^\mu = h P^\mu}$$

CI METTIAMO IN UN SR CON  $P^\mu = (P^0, 0, 0, \pm P^0)$   
COSÌ :

RICORDA CHE  $m=0 \Rightarrow v=c \Rightarrow E=P^0 \neq 0$   
E NON ESISTE UN SR DI RIPOSO

$$W^0 = \vec{P} \cdot \vec{J} = \pm P^0 J^3 \Rightarrow \underline{W^3 = W^0}$$
$$W^i = P^0 J^i \Rightarrow W^3 = \pm P^0 J^3$$

DUNQUE VEDIAMO:

$$W^0 |p\rangle = \underline{\pm P^0 J^3 |p\rangle} = W^3 |p\rangle = \underline{\pm P^0 h |p\rangle}$$

CONFRONTANDO



$$J^3 |p\rangle = h |p\rangle$$

QUINDE GENERALIZZANDO

$$h = \vec{J} \cdot \hat{P} \Rightarrow \boxed{h = \frac{\vec{J} \cdot \vec{P}}{|\vec{P}|}}$$



$$\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \delta \phi_a \right) - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \delta \phi_a$$

$$\phi_a(x) \rightarrow \phi_a(x) + \delta \phi_a(x)$$

$$\Rightarrow \delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi_a} \delta \phi_a + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \delta (\partial_\mu \phi_a) = \left( \frac{\partial \mathcal{L}}{\partial \phi_a} - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \right) \delta \phi_a + \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \delta \phi_a \right)$$

PRENDIAMO  $\delta \phi_a$  t.c.  $\delta S = 0 \Rightarrow \delta \mathcal{L} = \partial_\mu F^\mu$

SE ON-SHELL  $\Rightarrow \frac{\partial \mathcal{L}}{\partial \phi_a} - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} = 0 \Rightarrow \delta \mathcal{L} = \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \delta \phi_a \right) \equiv \partial_\mu F^\mu$

$$\Rightarrow J^\mu = \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_a]} \delta \phi_a - F^\mu \Rightarrow \partial_\mu J^\mu = 0 \Rightarrow J^\mu \text{ CONSERVATA}$$

(15)

$$\sigma^{0i} = \frac{i}{2} \left( \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix} - \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix} \right) = \frac{i}{2} \left( \begin{pmatrix} -\sigma^i & 0 \\ 0 & \sigma^i \end{pmatrix} - \begin{pmatrix} \sigma^i & 0 \\ 0 & -\sigma^i \end{pmatrix} \right) = i \begin{pmatrix} -\sigma^i & 0 \\ 0 & \sigma^i \end{pmatrix}$$

$$\begin{aligned} & \frac{i}{2} \begin{pmatrix} [\sigma^i, \sigma^j] & 0 \\ 0 & [\sigma^i, \sigma^j] \end{pmatrix} = -\varepsilon^{ijk} \begin{pmatrix} \sigma^k & 0 \\ 0 & \sigma^k \end{pmatrix} = \varepsilon^{ijk} \begin{pmatrix} \sigma^k & 0 \\ 0 & \sigma^k \end{pmatrix} \\ & \frac{i}{2} \left( \begin{pmatrix} -\sigma^i \sigma^j & 0 \\ 0 & -\sigma^i \sigma^j \end{pmatrix} - \begin{pmatrix} -\sigma^j \sigma^i & 0 \\ 0 & -\sigma^j \sigma^i \end{pmatrix} \right) \\ \sigma^{ij} &= \frac{i}{2} \left( \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^j \\ \sigma^j & 0 \end{pmatrix} - \begin{pmatrix} 0 & \sigma^j \\ \sigma^j & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix} \right) \end{aligned}$$

$$([\sigma^i, \sigma^j] = 2i \varepsilon^{ijk} \sigma^k)$$

(6)

SAPPIAMO

$$(\sigma^i)^{2m} = \mathbb{1}$$

$$(\sigma^i)^{2m+1} = \sigma^i$$

$$[\sigma^i, \sigma^j] = i \epsilon^{ijk} \sigma^k$$

IN RAPPRESENTAZIONE DI DIRAC

$$\sigma^0 = \begin{pmatrix} \mathbb{1} & 0 \\ 0 & -\mathbb{1} \end{pmatrix}$$

$$\sigma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}$$

$$[\sigma^0, \sigma^i] = 2 \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}$$

$$\sigma^0 \sigma^i = \frac{i}{2} [\sigma^0, \sigma^i] = i \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}$$

$$\left\{ \begin{aligned} \gamma^{\mu\nu} &= \frac{i}{2} \sigma^{\mu\nu} \\ \gamma^{0i} &= \gamma^i = \frac{i}{2} \sigma^{0i} \end{aligned} \right.$$

$$\gamma^{0i} = \gamma^i = \frac{i}{2} \sigma^{0i}$$

$$\begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}^2 = \mathbb{1}$$

$$\begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}^3 = \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}$$

$$\vec{\gamma} = \gamma \hat{n} = (\omega^0, \omega^1, \omega^2, \omega^3)$$

IN GENERALE DATA UNA MATRICE  $A$  E.C.  $A^{2m} = \mathbb{1}$   $A^{2m+1} = A$

$$\Rightarrow e^{\alpha A} \sim \mathbb{1} + (\alpha A) + \frac{1}{2!} (\alpha)^2 \mathbb{1} + \frac{1}{3!} (\alpha)^3 A + \dots$$

$$= \cosh(\alpha) \mathbb{1} + \sinh(\alpha) A$$

NOI

PRENDIAMO

$$A = \begin{pmatrix} 0 & \sigma^i \\ \sigma^i & 0 \end{pmatrix}$$



L'ESPONENTE DI  $S(\lambda)$  È:

$$-\frac{i}{4} \omega_{0z} \sigma^{0z} = -\frac{i}{4} \eta_z \cdot \begin{pmatrix} 0 & \sigma^z \\ \sigma^z & 0 \end{pmatrix} = \frac{\eta}{4} \begin{pmatrix} 0 & \vec{\sigma} \cdot \hat{n} \\ \vec{\sigma} \cdot \hat{n} & 0 \end{pmatrix} \equiv \frac{\eta}{4} A$$

QUINDI

$$\exp \left\{ -\frac{i}{4} \omega_{0z} \sigma^{0z} \right\} = 1 \cosh\left(\frac{\eta}{4}\right) + \cancel{\text{term}} \sinh\left(\frac{\eta}{4}\right) A$$
$$= \begin{pmatrix} \cosh(\frac{\eta}{4}) & (\vec{\sigma} \cdot \hat{n}) \sinh(\frac{\eta}{4}) \\ (\vec{\sigma} \cdot \hat{n}) \sinh(\frac{\eta}{4}) & \cosh(\frac{\eta}{4}) \end{pmatrix}$$

IL FATTORE 4 È CONVENZIONE,  
DEFINENDO

$$\vec{\eta} = \frac{1}{2} (\omega^0, \omega^1, \omega^2, \omega^3)$$

SI SAREBBE OTTENUTO COME  
ARGOMENTO DI  $\cosh$  E  $\sinh$   
UN FATTORE  $\eta/2$

OPPURE È ANCHE COMUNE

$$\sigma^{\mu\nu} = \frac{i}{4} [\sigma^\mu, \sigma^\nu]$$

$$\not{p}\not{p} = \gamma^\mu p_\mu \gamma^\nu p_\nu = \underbrace{\gamma^\mu \gamma^\nu}_{\text{}} p_\mu p_\nu$$

$$\gamma^\mu \gamma^\nu = -\gamma^\nu \gamma^\mu + \{\gamma^\mu, \gamma^\nu\}$$

$$= -\gamma^\nu \gamma^\mu + 2g^{\mu\nu}$$

$$= -(\gamma^\mu \gamma^\nu + [\gamma^\nu, \gamma^\mu]) + 2g^{\mu\nu}$$

$$= -\gamma^\mu \gamma^\nu + 2i\sigma^{\nu\mu} + 2g^{\mu\nu}$$

$$\Rightarrow 2\gamma^\mu \gamma^\nu = \underbrace{2i\sigma^{\nu\mu}}_{\text{ANTISYM.}} + \underbrace{2g^{\mu\nu}}_{\text{SYM.}}$$

$$\Rightarrow \not{p}\not{p} = [i\sigma^{\nu\mu} + g^{\mu\nu}] p_\mu p_\nu = p_\mu p^\mu = p^2$$

$$\langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle = \int \frac{d^3 p}{(2\pi)^3 2E_p} e^{-ip(x-y)} \quad (p, \theta, \varphi)$$

$$x^0 - y^0 > 0, \quad \vec{x} = \vec{y} \Rightarrow x^0 - y^0 = x^0 - y^0 > 0 \Rightarrow \langle 0 | \phi(x) \phi(y) | 0 \rangle = \int \frac{d^3 p}{(2\pi)^3 2E_p} e^{-ip(x-y)}$$

$$\begin{aligned} E_p &= \sqrt{\vec{p}^2 + m^2} \\ E_p^2 - m^2 &= \vec{p}^2 \\ E_p dE_p &= d^3 p \\ dp &= \frac{E_p}{\sqrt{E_p^2 - m^2}} dE_p \end{aligned}$$

$$= \frac{1}{(4\pi^2)} \int_{-m}^{\infty} \frac{dE \sqrt{E^2 - m^2}}{E_p} e^{-iEt} \xrightarrow{t \rightarrow \infty} 0$$



$$\langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle = \int \frac{d^3 p}{(2\pi)^3 2E_p} e^{-ip(x-y)} \quad (p, \theta, \varphi)$$

$$= \frac{1}{(2\pi)^3} \int \frac{p^2 dp}{2E_p} e^{-ip \cdot (\vec{x} - \vec{y})} \int m d\theta \int d\varphi$$

$$x^0 - y^0 = \Delta(x-y) \underset{t \rightarrow \infty}{\sim} e^{-im|t|}$$

$$\vec{x} - \vec{y} =$$

$$= \frac{2\pi}{(2\pi)^3} \int \frac{p^2 dp}{E_p} e^{-ipz \cos \theta} m d\theta$$

$$= \frac{1}{ipz} \frac{d}{d\theta} (e^{-ipz \cos \theta})$$

$$= \frac{1}{(2\pi)^2} \int \frac{p^2 dp}{E_p} \frac{d\theta}{ipz} \frac{d}{d\theta} (e^{-ipz \cos \theta})$$

(21)

$$|\phi(x)\phi^\dagger(y)|0\rangle = \frac{1}{(2\pi)^2} \int \frac{dp}{E_p} \frac{p^2}{ip^2} \underbrace{\int_0^{2\pi} d\theta \frac{d}{d\theta}}_{\text{}} (e^{-ip^2 \cos \theta})$$

$$\gamma^0 : \Delta(x-y) \underset{t \rightarrow \infty}{\sim} e^{-imt}$$

$$= \frac{1}{(2\pi)^2} \int_0^\infty \frac{dp}{E_p} \frac{p}{ip^2} (e^{ip^2} - e^{-ip^2})$$

$$= \frac{-1}{(2\pi)^2} \int_0^\infty \frac{p}{\sqrt{p^2+m^2}} (e^{ip^2} - e^{-ip^2}) \quad p = |p|$$

$$= \frac{-1}{(2\pi)^2} \left\{ \int_0^\infty \frac{p}{\sqrt{p^2+m^2}} e^{ip^2} - \int_0^\infty \frac{p}{\sqrt{p^2+m^2}} e^{-ip^2} \right\}$$

$$p \rightarrow -p \quad \Rightarrow \int_0^\infty \frac{-p}{\sqrt{p^2+m^2}} e^{-ip^2} = - \int_{-\infty}^0 \frac{p}{\sqrt{p^2+m^2}} e^{-ip^2}$$

$$\langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle = \frac{-i}{(2\pi)^2 z} \left\{ \int_0^{+\infty} dp \frac{p}{\sqrt{p^2 + m^2}} e^{-ipz} + \int_{-\infty}^0 dp \frac{p}{\sqrt{p^2 + m^2}} e^{ipz} \right\}$$

$$= \frac{-i}{(2\pi)^2 z} \int_{-\infty}^{+\infty} dp \frac{p}{\sqrt{p^2 + m^2}} e^{-ipz} \rightarrow \text{poles at } \pm im$$

$$s = -ip \quad \left[ \int_m^{\infty} ds \frac{s}{\sqrt{m^2 - s^2}} e^{-sz} \sim e^{-mz} \right] \quad \begin{matrix} x^0 = y^0 \\ \vec{x} - \vec{y} = x^1 - y^1 \end{matrix}$$

(23)