

USANDO  $H^{uv} = i(x^u \partial^v - x^v \partial^u)$  VALUTAMO  $[H^{uv}, H^{rs}]$ .  
 PER NON PERDERE PEZZI SERVIAMOCI DI UNA FUNZIONE  $f = f(x^u)$ ;

$$\begin{aligned}
 H^{uv} H^{rs} f &= -i(x^u \partial^v - x^v \partial^u)(x^s \partial^r - x^r \partial^s) f \\
 &= -[x^u \partial^v(x^s \partial^r) - x^u \partial^v(x^r \partial^s) - x^v \partial^u(x^s \partial^r) + x^v \partial^u(x^r \partial^s)] f \\
 &= -[x^u(g^{vs} \partial^r + x^s \partial^v \partial^r) - x^u(g^{ur} \partial^s + x^r \partial^u \partial^s) + - \\
 &\quad - x^v(g^{us} \partial^r + x^s \partial^u \partial^r) + x^v(g^{vr} \partial^s + x^r \partial^v \partial^s)] f
 \end{aligned}$$

$\begin{aligned} \partial^v &= g^{v\alpha} \partial_\alpha \\ \partial^v x^s &= g^{v\alpha} \partial_\alpha x^s \\ &= g^{v\alpha} \delta_{\alpha}^s \\ &= g^{vs} \end{aligned}$  IN QUESTO MODO:

$$\begin{aligned}
 \boxed{[H^{uv}, H^{rs}]f} &= [H^{uv} H^{rs} - H^{rs} H^{uv}]f = \\
 &= -[x^u g^{vs} \partial^r + x^u x^s \partial^v \partial^r - \cancel{x^u g^{vu} \partial^s} - \cancel{x^u x^s \partial^v \partial^s} - \cancel{x^v g^{us} \partial^r} - \cancel{x^v x^s \partial^u \partial^r} + \\
 &\quad + \cancel{x^v g^{vu} \partial^s} + \cancel{x^v x^s \partial^u \partial^s}] f + \boxed{[x^s g^{ur} \partial^v + x^s x^r \partial^u \partial^v - \cancel{x^s g^{vu} \partial^r} - \cancel{x^s x^v \partial^u \partial^r} - \\
 &\quad - \cancel{x^r g^{us} \partial^v} - \cancel{x^r x^u \partial^s \partial^v} + \cancel{x^r g^{vu} \partial^s} + \cancel{x^r x^v \partial^u \partial^s}] f}
 \end{aligned}$$

VISTA LA  
SIMMETRIA  
 $\begin{aligned} &g^{uv} \\ &\uparrow \end{aligned}$

TERMINI  
RACCOLIERE:

SOTTOLINEATI A ZIG-ZAG SI CANCELLANO, MENTRE GLI ALTRI LI POSSIAMO

$$\boxed{[H^{uv}, H^{rs}] = -[g^{vs}(x^u \partial^r - x^r \partial^u) + g^{ur}(x^s \partial^v - x^v \partial^s) + \dots]} \quad \textcircled{1}$$

$$\dots + \underline{g^{\mu\beta}(x^\sigma \partial^\nu - x^\nu \partial^\sigma)} + \underline{g^{\nu\sigma}(x^\beta \partial^\mu - x^\mu \partial^\beta)} \Big]$$

POSSIAMO RICONOSCERE :  $-\frac{1}{2}M^{\mu\nu} = (x^\mu \partial^\nu - x^\nu \partial^\mu)$

$$[M^{\mu\nu}, M^{\beta\sigma}] = -i [g^{\nu\beta}M^{\mu\sigma} + g^{\nu\sigma}M^{\beta\mu} + g^{\mu\beta}M^{\nu\sigma} + g^{\mu\sigma}M^{\nu\beta}]$$

POSSIAMO ANCHE SISTEMARE I SEGNI USANDO LA SIMMETRIA DI  $g^{\mu\nu}$  E L'ANTISIMMETRIA DI  $M^{\mu\nu}$ . LA STRUTTURA DEGLI INDICI DEI TESTI È :

$$[12, 34] = -\frac{1}{2} (14, 23) + (23, 14) - (13, 24) - (24, 13)$$

INDICI  $M$   
INDICI  $g$

SISTEMANDO LA NOSTRA SCRITTURA PER OTTENERE :

$$[M^{\mu\nu}, M^{\beta\sigma}] = -\frac{1}{2} [(\mu\sigma, \nu\beta) + (\nu\beta, \mu\sigma) - (\mu\beta, \nu\sigma) - (\nu\sigma, \mu\beta)]$$

NOI ABBIANO

$$= -\frac{1}{2} [g^{\mu\sigma}M^{\nu\beta} + g^{\nu\beta}M^{\mu\sigma} - g^{\mu\beta}M^{\nu\sigma} - g^{\nu\sigma}M^{\mu\beta}]$$

AVENDO

$$N_{\dot{z}} = \frac{1}{2} (J_{\dot{z}} + \dot{z} K_{\dot{z}})$$

$$N_{\dot{z}}^+ = \frac{1}{2} (J_{\dot{z}} - \dot{z} K_{\dot{z}})$$

E' FACILE

VERIFICARE

$$\begin{aligned}
 [N_{\dot{z}}, N_{\dot{z}}^+] &= \left[ \frac{1}{2} (J_{\dot{z}} + \dot{z} K_{\dot{z}}), \frac{1}{2} (J_{\dot{z}} - \dot{z} K_{\dot{z}}) \right] \\
 &= \frac{1}{4} \left( [J_{\dot{z}}, J_{\dot{z}}] - \dot{z} [J_{\dot{z}}, K_{\dot{z}}] + \dot{z} [K_{\dot{z}}, J_{\dot{z}}] + [K_{\dot{z}}, K_{\dot{z}}] \right) \\
 &= \frac{1}{4} \left( \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K + \dot{z} \epsilon_{\dot{z} \dot{z} K} K_K + \dot{z} \left( -\underbrace{\dot{z} \epsilon_{\dot{z} \dot{z} K} K_K}_{-\epsilon_{\dot{z} \dot{z} K}} \right) - \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K \right) \\
 &= \frac{1}{4} \left( \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K + \epsilon_{\dot{z} \dot{z} K} K_K - \epsilon_{\dot{z} \dot{z} K} K_K - \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K \right) \\
 &= 0
 \end{aligned}$$

$$\begin{aligned}
 [N_{\dot{z}}, N_{\dot{z}}] &= \frac{1}{4} [J_{\dot{z}} + \dot{z} K_{\dot{z}}, J_{\dot{z}} + \dot{z} K_{\dot{z}}] = \frac{1}{4} \left( [J_{\dot{z}}, J_{\dot{z}}] + \dot{z} [J_{\dot{z}}, K_{\dot{z}}] + \dot{z} [K_{\dot{z}}, J_{\dot{z}}] - [K_{\dot{z}}, K_{\dot{z}}] \right) \\
 &= \frac{1}{4} \left( \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K + \dot{z} \left( \dot{z} \epsilon_{\dot{z} \dot{z} K} K_K \right) + \dot{z} \left( -\underbrace{\dot{z} \epsilon_{\dot{z} \dot{z} K} K_K}_{-\epsilon_{\dot{z} \dot{z} K}} \right) + \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K \right) \\
 &= \frac{1}{2} \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K - \frac{1}{2} \dot{z} \epsilon_{\dot{z} \dot{z} K} K_K \\
 &= \dot{z} \epsilon_{\dot{z} \dot{z} K} \left( \frac{1}{2} (J_K + \dot{z} K_K) \right) = \dot{z} \epsilon_{\dot{z} \dot{z} K} N_K
 \end{aligned}$$

VALGONO:

$$[J_{\dot{z}}, J_{\dot{z}}] = \dot{z} \epsilon_{\dot{z} \dot{z} K} J_K$$

$$[K_{\dot{z}}, K_{\dot{z}}] = -\dot{z} \epsilon_{\dot{z} \dot{z} K} J_K$$

$$[J_{\dot{z}}, K_{\dot{z}}] = \dot{z} \epsilon_{\dot{z} \dot{z} K} K_K$$

$$\begin{aligned}
 [A, B] &= AB - BA \\
 [B, A] &= BA - AB \\
 &= -[A, B]
 \end{aligned}$$

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$$\begin{aligned}
 & [N_z^+, N_j^+] \cdot \frac{1}{4} [J_z - \dot{\alpha} K_z, J_j - \dot{\alpha} K_j] \\
 &= \frac{1}{4} \left( [J_z, J_j] - \dot{\alpha} [J_z, K_j] - \dot{\alpha} [K_z, J_j] - [K_z, K_j] \right) \\
 &\quad - [J_j, K_z] = -\dot{\alpha} \varepsilon_{ijk} K_k = \dot{\alpha} \varepsilon_{ijk} K_k \\
 &= \frac{1}{4} \left( \dot{\alpha} \varepsilon_{ijk} J_k - \dot{\alpha} (\dot{\alpha} \varepsilon_{ijk} K_k) - \dot{\alpha} (\dot{\alpha} \varepsilon_{ijk} K_k) + \dot{\alpha} \varepsilon_{ijk} K_k \right) \\
 &= \frac{1}{2} \dot{\alpha} \varepsilon_{ijk} J_k + \frac{1}{2} \dot{\alpha} \varepsilon_{ijk} K_k \\
 &= \dot{\alpha} \varepsilon_{ijk} \frac{1}{2} (J_k - \dot{\alpha} K_k) \\
 &= \dot{\alpha} \varepsilon_{ijk} N_k^+
 \end{aligned}$$

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CONOSCO

$$(M^{1\sigma})^M_{\sim} = -\frac{1}{2} (g^{M\sigma} \delta^1_{\sim} - g^{M\sigma} \delta^0_{\sim})$$

POSSIAMO DETERMINARE LE 6 MATRICI  $M^{i\sigma}$ .

$$(M^{01})^M_{\sim} = -\frac{1}{2} (g^{M0} \delta^1_{\sim} - g^{M1} \delta^0_{\sim}) = -\frac{1}{2} \left[ \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} (0100) - \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} (01000) \right] = -\frac{1}{2} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{02})^M_{\sim} = -\frac{1}{2} (g^{M0} \delta^2_{\sim} - g^{M2} \delta^0_{\sim}) = -\frac{1}{2} \left[ \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} (0010) - \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix} (1000) \right] = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{03})^M_{\sim} = -\frac{1}{2} (g^{M0} \delta^3_{\sim} - g^{M3} \delta^0_{\sim}) = -\frac{1}{2} \left[ \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} (0001) - \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} (1000) \right] = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{12})^M_{\sim} = -\frac{1}{2} (g^{M1} \delta^2_{\sim} - g^{M2} \delta^1_{\sim}) = -\frac{1}{2} \left[ \begin{pmatrix} 0 \\ -1 \\ 0 \\ 0 \end{pmatrix} (0010) - \begin{pmatrix} 0 \\ 0 \\ -1 \\ 0 \end{pmatrix} (0100) \right] = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$(M^{13})^M_{\sim} = -\frac{1}{2} (g^{M1} \delta^3_{\sim} - g^{M3} \delta^1_{\sim}) = -\frac{1}{2} \left[ \begin{pmatrix} 0 \\ -1 \\ 0 \\ 0 \end{pmatrix} (0001) - \begin{pmatrix} 0 \\ 0 \\ 0 \\ -1 \end{pmatrix} (0100) \right] = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

$$(M^{13})^M_{\sim} = -\frac{1}{2} (g^{M1} \delta^3_{\sim} - g^{M3} \delta^1_{\sim}) = -\frac{1}{2} \left[ \begin{pmatrix} 0 \\ -1 \\ 0 \\ 0 \end{pmatrix} (0001) - \begin{pmatrix} 0 \\ 0 \\ 0 \\ -1 \end{pmatrix} (0100) \right] = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

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POSSIAMO

VERIFICARE

CHE

UGUALE A  $J_z = \frac{1}{2} \varepsilon_{ijk} M_{ik}$  POICHÉ ABBASSANDO  $z$   
INDICI SPAZIALI ESCE "(-)(-) = +"

$$J_z = \frac{1}{2} \varepsilon_{ijk} M_{ik}$$

$$K^z = M^{0z}$$

È IMMEDIATO

$$K^x = K_x = M^{01} = -\frac{1}{2} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$K^y = K_y = M^{02} = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

VEDI  
ILIOPOULOS  
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$$K^3 = K_z = M^{03} = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

$$J^1 = J^x = \frac{1}{2} \varepsilon^{ijk} M_{ik} = \frac{1}{2} \left( \underbrace{\varepsilon^{123} M^{23}}_{(-\varepsilon^{123})(-M^{23})} + \underbrace{\varepsilon^{132} M^{32}}_{+1} \right) = \varepsilon^{123} M^{23} = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & +1 & 0 \end{pmatrix}$$

$$J^2 = J^y = \frac{1}{2} \varepsilon^{2jk} M_{ik} = \frac{1}{2} \left( \underbrace{\varepsilon^{213} M^{13}}_{-\varepsilon^{123}} + \underbrace{\varepsilon^{231} M^{31}}_{\varepsilon^{123}(-M^{13})} \right) = - (M^{13}) = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & +1 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$J^3 = J^z = \frac{1}{2} \varepsilon^{3jk} M_{ik} = \frac{1}{2} \left( \underbrace{\varepsilon^{312} M^{12}}_{\varepsilon^{123} M^{12}} + \underbrace{\varepsilon^{321} M^{21}}_{(-\varepsilon^{123})(-M^{12})} \right) = M^{12} = -\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

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VEDIAMO I MOMENTI ANGOLARI  $J_1$  E I BOOST  $K_1$  COSA DIVENTANO PER TRASFORMAZIONI DI PARITÀ.

VEDIAMO AD ESEMPIO  $J_1$ :

$$\Lambda_p = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

VEI § 3.7.2  
DI SCHWICHTENBERG

$$J_1' = \Lambda_p J_1 \Lambda_p^T = \Lambda_p J_1 \Lambda_p$$

$$= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}^T$$

$$= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} = J_1$$

ANALOGHI CONTI PER  $J_2$  E  $J_3$ . PER IL BOOST  $K_1$ :

$$K_1' = \Lambda_p K_1 \Lambda_p^T = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}^T = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$= \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} = -K_1$$

CONTI ANALOGHI PER  $K_2$ ,  $K_3$ .

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CON CONTI ANALOGHI SI VEDA PER L'INVERSIONE TEMPORALE:

$$J_{\hat{z}} \xrightarrow{T} J'_{\hat{z}} = \Lambda_T J_{\hat{z}} \Lambda_T^T = J_{\hat{z}}$$

$$K_{\hat{z}} \xrightarrow{T} K'_{\hat{z}} = \Lambda_T K_{\hat{z}} \Lambda_T^T = -K_{\hat{z}}$$

QUINDI COMPLESSIVAMENTE SCRIVIAMO

$$J_{\hat{z}} \xrightarrow{P} J_{\hat{z}}$$

$$K_{\hat{z}} \xrightarrow{P} -K_{\hat{z}}$$

$$J_{\hat{z}} \xrightarrow{T} J_{\hat{z}}$$

$$K_{\hat{z}} \xrightarrow{T} -K_{\hat{z}}$$

$$\Lambda_{L,R} = e^{+\frac{i}{2}\sigma^z(\omega_{\pm} \mp i\eta_{\pm})}$$

$$\Lambda_{L,R}^* = e^{-\frac{i}{2}\sigma^z(\omega_{\pm} \pm i\eta_{\pm})}$$

$$\Lambda_{L,R}^T = e^{+\frac{i}{2}\sigma^z(\omega_{\pm} \pm i\eta_{\pm})}$$

$$\Lambda_{L,R}^{-1} = e^{-\frac{i}{2}\sigma^z(\omega_{\pm} \mp i\eta_{\pm})}$$

$$\Lambda_{L,R}^+ = e^{\frac{i}{2}\sigma^z(\omega_{\pm} \mp i\eta_{\pm})}$$

$$\begin{cases} (\sigma^2)^* = -\sigma^2 & ; (\sigma^2)^T = -\sigma^2 \\ (\sigma^2)^* = \sigma^2 & ; (\sigma^2)^+ = \sigma^2 \end{cases}$$

$$\sigma^2 \sigma^z \sigma^2 = -\sigma^z \Rightarrow -\sigma^z = \sigma^2 \sigma^z \sigma^2$$

$$\textcircled{1} \quad \Lambda_{L,R}^* \xrightarrow{\quad} \Lambda_{L,R}^+ = e^{\frac{i}{2}\sigma^z(\omega_{\pm} \pm i\eta_{\pm})}$$

$$\Rightarrow \boxed{\Lambda_{L,R}^+ = \Lambda_{R,L}^{-1}}$$

$$\textcircled{2} \quad \boxed{\sigma^2 \Lambda_{L,R}^* \sigma^2 = e^{+\frac{i}{2}\sigma^z(\omega_{\pm} \pm i\eta_{\pm})} = \Lambda_{R,L}} \quad \textcircled{3} \quad \boxed{\Lambda_{L,R}^* = \sigma^2 \Lambda_{R,L} \sigma^2}$$

$$\textcircled{4} \quad \textcircled{3} \xrightarrow{+ = * + T} (\Lambda_{L,R}^*)^+ = (\sigma^2 \Lambda_{R,L} \sigma^2)^+ \Rightarrow \Lambda_{L,R}^T = \sigma^2 \Lambda_{R,L}^+ \sigma^2 \xrightarrow{\textcircled{1}} \boxed{\Lambda_{L,R}^T = \sigma^2 \Lambda_{L,R}^{-1} \sigma^2}$$

$$\textcircled{5} \quad \boxed{\Lambda_L^T \sigma^2 \Lambda_L = \Lambda_L^T \sigma^2 \Lambda_L \sigma^2 \sigma^2 = \Lambda_L^T \Lambda_R^* \sigma^2 \stackrel{\textcircled{3}}{=} \sigma^2 \Lambda_L^{-1} \sigma^2 \Lambda_R^* \sigma^2 = \sigma^2 \Lambda_L^{-1} \Lambda_L = \boxed{\sigma^2}}$$

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CONTO  
GUARDO

GENERALICO PER  $\phi(x^\mu)$ ,  $\mathcal{L}(\phi, \partial_\mu \phi)$

$$\Delta S = \int d^4x \times \Delta \mathcal{L} + \int \delta(\partial^\mu x) \mathcal{L}$$

so  $\delta(\partial^\mu x) = \partial_\mu \delta x^\mu d^\mu x$   
PER UNA TRASFORMAZIONE  
DI COORDINATE (VEDI  
NOTE)

CALCOLO

$$\begin{aligned}\Delta \mathcal{L} &= \delta \mathcal{L} + \delta x^\mu \partial_\mu \mathcal{L} \\ &= \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta(\partial_\mu \phi) + \delta x^\mu \partial_\mu \mathcal{L} \\ &= \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \partial_\mu(\delta \phi) + \delta x^\mu \partial_\mu \mathcal{L} \\ &= \frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi \right) - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi + \delta x^\mu \partial_\mu \mathcal{L}\end{aligned}$$

MA NETTO ON-SHELL  $\Rightarrow$  EQ MOTTO = 0

$$\Rightarrow \Delta \mathcal{L} = \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi \right) + \delta x^\mu \partial_\mu \mathcal{L}$$

così

$$\Delta S = \int d^4x \times \underbrace{\left[ \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi \right) + \delta x^\mu \partial_\mu \mathcal{L} + \partial_\mu \delta x^\mu \mathcal{L} \right]}_{\partial_\mu(\delta x^\mu \mathcal{L})}$$

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$$\Rightarrow \Delta S = \int d^4x \underbrace{\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi + \delta x^\mu \mathcal{L} \right)}_{J^\mu} = 0$$

QUINDI ABBIAMO

$$J^\mu = \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta \phi + \delta x^\mu \mathcal{L}$$

SE DIFFERENZIO I CASI DI TRASFORMAZIONE DI COORDINATE:  
 $x^\mu \rightarrow x^\mu + \delta x^\mu \Rightarrow \Delta \phi = 0 \Rightarrow \delta \phi = - \delta x^\mu \partial_\mu \phi$

SOSTITUISCO

$$J^\mu = - \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \delta x^\nu \partial_\nu \phi + \delta x^\mu \mathcal{L} = \left( - \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \partial_\nu \phi + g^{\mu\nu} \mathcal{L} \right) \delta x^\nu = - T^\mu_\nu \delta x^\nu$$

TENSORE  
ENERGIA-IMPULSO

VEDIAMO ANCHE IL CASO GENERICO:

$$x^\mu \rightarrow x^\mu + \delta x^\mu, \quad \phi \rightarrow \phi + \alpha \Rightarrow \Delta \phi = \delta \phi + \delta x^\mu \partial_\mu \phi$$

SOSTITUENDO  ~~$\delta \phi = \Delta \phi - \delta x^\mu \partial_\mu \phi$~~

$$J^\mu = - T^\mu_\nu \delta x^\nu + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi]} \Delta \phi$$

(4)

$$\delta S = S(x) - S(x) = \int d^4x \left( \mathcal{L}'(x) - \mathcal{L}(x) \right) = \int d^4x \delta \mathcal{L}$$

$$d^4x' = (1 + \partial_\mu \delta x^\mu) d^4x$$

$$\Delta S = S'(x') - S(x) = \int d^4x' \mathcal{L}'(x') - \int d^4x \mathcal{L}(x) \quad \Rightarrow \quad \mathcal{L}'(x) = \Delta \mathcal{L} + \mathcal{L}(x)$$

$$= \int (1 + \partial_\mu \delta x^\mu) (\Delta \mathcal{L} + \mathcal{L}(x)) d^4x - \int d^4x \mathcal{L}(x)$$

$$\simeq \int d^4x \left( \Delta \mathcal{L} + \partial_\mu \delta x^\mu \mathcal{L} \right) = \int d^4x \left( \Delta \mathcal{L} + \underbrace{\delta \mathcal{L} + \delta x^\mu \partial_\mu \mathcal{L} + \mathcal{L} \partial_\mu \delta x^\mu}_{\partial_\mu (\delta x^\mu \mathcal{L})} \right)$$

$$\Delta \mathcal{L} = \delta \mathcal{L} + \delta x^\mu \partial_\mu \mathcal{L}$$

$$\Delta S = \delta S + \int d^4x \partial_\mu (\mathcal{L} \delta x^\mu) \quad (12)$$

DIMOSTRAZIONE

$$h = \frac{\vec{j} \cdot \vec{p}}{|\vec{p}|}$$

PART. A VO DAL SAPERE:

$$W^{\mu} = \frac{1}{2} \epsilon^{\mu\nu\gamma\sigma} P_{\nu} M_{\gamma\sigma}$$

$$\left\{ \begin{array}{l} W^0 = \frac{1}{2} \epsilon^{0\hat{i}jk} P_{\hat{i}} M_{jk} = P_{\hat{i}} J^{\hat{i}} = \vec{P} \cdot \vec{J} \\ \vec{W} = P_0 \vec{J} - \vec{P} \times \vec{K} \end{array} \right.$$

PRENDIAMO  $m=0 \Rightarrow P^{\mu} = (0, \vec{p})$

$$\Rightarrow \begin{cases} P^{\mu} P^{\mu} = 0 \\ W_{\mu} W^{\mu} = 0 \\ W_{\mu} P^{\mu} = 0 \end{cases} \Rightarrow \boxed{W^{\mu} = h P^{\mu}}$$

Ci mettiamo in un SR con  $P^{\mu} = (P^0, 0, 0, \pm P^0)$

Così:

RICORDA CHE  $m=0 \Rightarrow v=c \Rightarrow E=P^0 \neq 0$   
E NON ESISTE UN SR DI RIPOSO

$$W^0 = \vec{P} \cdot \vec{J} = \pm P^0 J^3 \Rightarrow \underline{W^3 = W^0}$$

$$W^1 = P^0 J^1 \Rightarrow W^3 = \pm P^0 J^3$$

DUNQUE VEDIAMO:

$$W^0 |p\rangle = \underline{\pm P^0 J^3 |p\rangle} = W^3 |p\rangle = \underline{\pm P^0 h |p\rangle}$$

CONFRONTANDO  
➡

$$\vec{J}^3 |p\rangle = h |p\rangle$$

QUNQUE GENERALIZZANDO

$$h = \vec{J} \cdot \hat{P} \Rightarrow$$

$$h = \frac{\vec{J} \cdot \vec{P}}{|\vec{P}|}$$

$$\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \delta \phi_e \right) - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \delta \phi_e$$

$$\phi_e^{(x)} \rightarrow \phi_e^{(x)} + \delta \phi_e^{(x)}$$

↑

$$\Rightarrow \delta \mathcal{L} = \frac{\partial \mathcal{L}}{\partial \phi_e} \delta \phi_e + \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \delta (\partial_\mu \phi_e) = \left( \frac{\partial \mathcal{L}}{\partial \phi_e} - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \right) \delta \phi_e + \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \delta \phi_e \right)$$

PRENDIAMO  $\delta \phi_e$  t.c.  $\delta S = 0 \Rightarrow \delta \mathcal{L} = \partial_\mu F^\mu$

SE ON-SHELL  $\Rightarrow \frac{\partial \mathcal{L}}{\partial \phi_e} - \partial_\mu \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} = 0 \Rightarrow \delta \mathcal{L} = \partial_\mu \left( \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \delta \phi_e \right) = \partial_\mu F^\mu$

$$\Rightarrow \bar{J}^\mu = \frac{\partial \mathcal{L}}{\partial [\partial_\mu \phi_e]} \delta \phi_e - F^\mu \Rightarrow \partial_\mu \bar{J}^\mu = 0 \Rightarrow \bar{J}^\mu \text{ CONSERVATA}$$

(15)

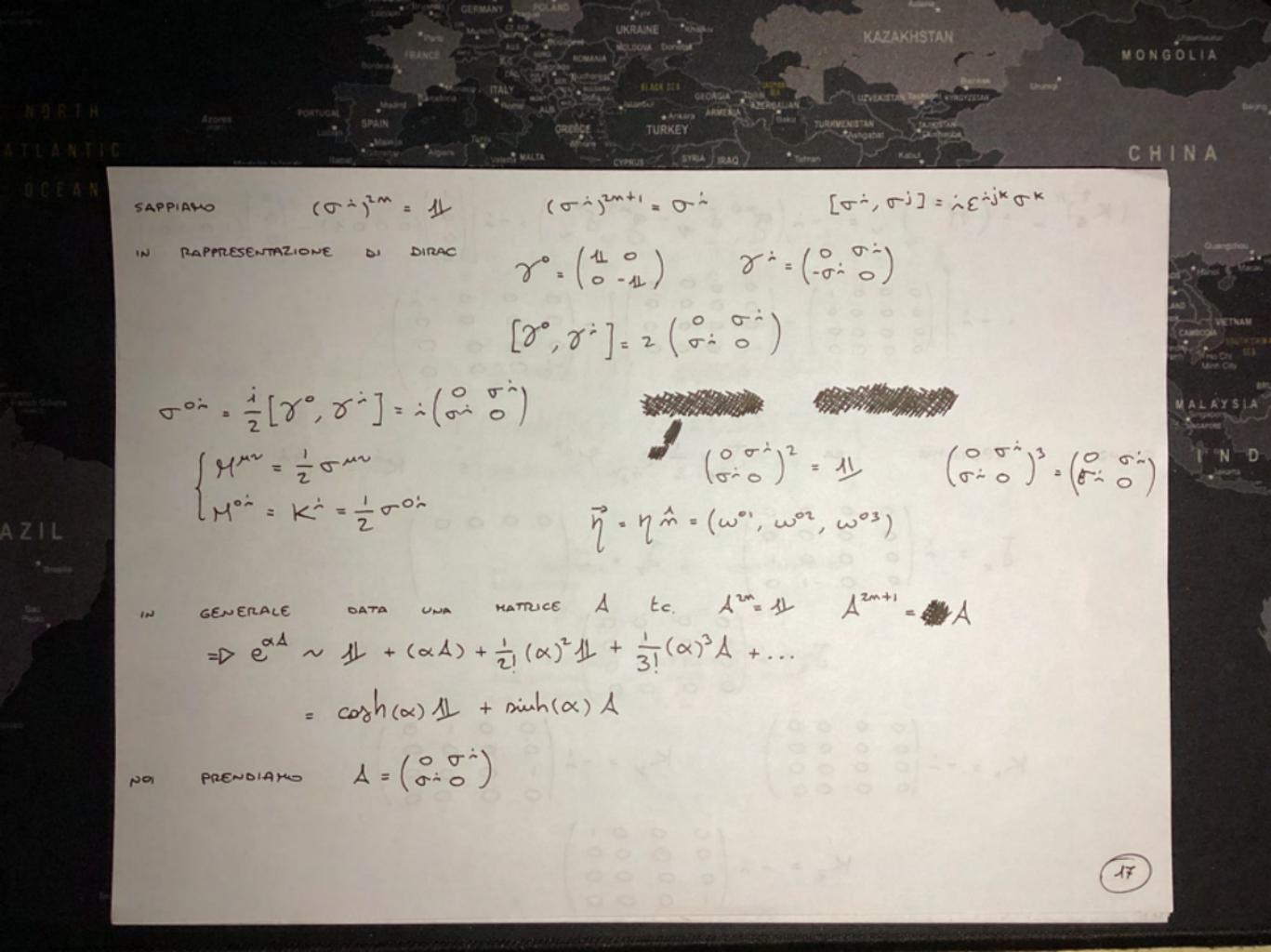
$$\sigma^{\pm} = \frac{1}{2} \left( \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^{\pm} \\ \sigma^{\pm} & 0 \end{pmatrix} - \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right) = \frac{1}{2} \left( \begin{pmatrix} -\sigma^{\pm} & 0 \\ 0 & \sigma^{\pm} \end{pmatrix} - \begin{pmatrix} 0 & 0 \\ 0 & -\sigma^{\pm} \end{pmatrix} \right) = \begin{pmatrix} -\sigma^{\pm} & 0 \\ 0 & \sigma^{\pm} \end{pmatrix}$$

$$\frac{1}{2} \begin{pmatrix} [\sigma^i \sigma^j] & 0 \\ 0 & [\sigma^i, \sigma^j] \end{pmatrix} = -\epsilon^{ijk} \begin{pmatrix} \sigma^k & 0 \\ 0 & \sigma^k \end{pmatrix} = \epsilon^{ijk} \begin{pmatrix} 0 & 0 \\ 0 & \sigma^k \end{pmatrix}$$

$$\sigma^{\pm} = \frac{1}{2} \left( \begin{pmatrix} -\sigma^{\pm} & 0 \\ 0 & -\sigma^{\pm} \end{pmatrix} - \begin{pmatrix} -\sigma^{\pm} & 0 \\ 0 & -\sigma^{\pm} \end{pmatrix} \right)$$

$$\sigma^{\pm} = \frac{1}{2} \left( \begin{pmatrix} 0 & \sigma^{\pm} \\ -\sigma^{\pm} & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^{\pm} \\ \sigma^{\pm} & 0 \end{pmatrix} - \begin{pmatrix} 0 & 0 \\ -\sigma^{\pm} & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma^{\pm} \\ \sigma^{\pm} & 0 \end{pmatrix} \right) \quad \left( [\sigma^i, \sigma^j] = \epsilon^{ijk} \sigma^k \right)$$

(16)



BC L'ESPONENTE DI  $S(\lambda)$  È:

$$-\frac{i}{4}\omega_0\hat{\sigma}^0\hat{n} = -\frac{i}{4}\eta_i\hat{n}\begin{pmatrix} 0 & \hat{\sigma}^z \\ \hat{\sigma}^z & 0 \end{pmatrix} = \frac{\eta}{4}\begin{pmatrix} 0 & \hat{\sigma}^z \cdot \hat{n} \\ \hat{\sigma}^z \cdot \hat{n} & 0 \end{pmatrix} = \frac{\eta}{4}A$$

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$$\exp\left\{-\frac{i}{4}\omega_0\sigma^z\right\} = \begin{pmatrix} 1 & \cancel{\text{cosh}}\left(\frac{n}{4}\right) + \cancel{\text{sinh}}\left(\frac{n}{4}\right) \\ \cancel{\text{cosh}}\left(\frac{n}{4}\right) & (\vec{\sigma} \cdot \vec{n}) \sinh\left(\frac{n}{4}\right) \\ (\vec{\sigma} \cdot \vec{n}) \sinh\left(\frac{n}{4}\right) & \text{cosh}\left(\frac{n}{4}\right) \end{pmatrix}$$

IL FATTORE 4 È CONVENZIONE,  
DEFINENDO

$$\vec{\eta} = \frac{1}{2}(\omega^1, \omega^2, \omega^3)$$

SI SAREBBE OTTENUTO COME ARGOMENTO DI  $\cosh^{-1} x$  UN FATTORE  $\frac{1}{2}$

OPPURE È ANCHE COMUNE

$$\sigma^{\mu\nu} = \frac{i}{4} [\gamma^\mu, \gamma^\nu]$$

$$\begin{aligned}
 p \bar{p} &= \gamma^\mu p_\mu \gamma^\nu p_\nu = \underbrace{\gamma^\mu \gamma^\nu}_{\gamma^\mu \gamma^\nu = -\gamma^\nu \gamma^\mu + \{\gamma^\mu, \gamma^\nu\}} p_\mu p_\nu \\
 &= -\gamma^\nu \gamma^\mu + 2g^{\mu\nu} p_\mu p_\nu \\
 &= -(\gamma^\mu \gamma^\nu + [\gamma^\nu, \gamma^\mu]) + 2g^{\mu\nu} p_\mu p_\nu \\
 &= -\gamma^\mu \gamma^\nu + 2i\sigma^{\mu\nu} + 2g^{\mu\nu} p_\mu p_\nu \\
 \Rightarrow 2\gamma^\mu \gamma^\nu &= 2i\sigma^{\mu\nu} + 2g^{\mu\nu} p_\mu p_\nu
 \end{aligned}$$

$$\Rightarrow p \bar{p} = [i\sigma^{\mu\nu} + g^{\mu\nu}] p_\mu p_\nu = p_\mu p^\mu = p^2$$

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 $\langle \phi | \Gamma \phi \rangle$ 

$$2) (G(x, y) \cdot p)(x-y) = \int \frac{d^3 p}{(2\pi)^3 2E_p} e^{-i\vec{p} \cdot (\vec{x}-\vec{y})} (p, \theta, \varphi)$$

$$\langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle = \int \frac{d^3 p}{(2\pi)^3 2E_p} e^{-i\vec{p} \cdot (\vec{x}-\vec{y})} (p, \theta, \varphi)$$

$$x^m - y^m > 0 \quad ; \quad \vec{x} = \vec{y} \Rightarrow x^m - y^m = x^0 - y^0 > 0 \Rightarrow$$

$$\begin{aligned} \langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle &= \int \frac{d^3 p}{(2\pi)^3 2E_p} e^{-i\vec{p} \cdot (\vec{x}-\vec{y})} (p, \theta, \varphi) \\ &= \frac{1}{2\pi^2} \int \frac{dp}{2\sqrt{p^2 + m^2}} e^{-i\vec{p} \cdot (\vec{x}-\vec{y})} \\ &= \frac{1}{(4\pi^2)} \int \frac{dE (E^2 - m^2)^{1/2}}{\sqrt{E^2 - m^2} E_p} e^{-iE \frac{(\vec{x}-\vec{y})}{E}} \\ &= \frac{1}{4\pi^2} \int_{-m}^{\infty} dE \sqrt{E^2 - m^2} e^{iEt} \underset{t \rightarrow \infty}{\propto} e^{iEt} \end{aligned}$$

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$$\langle \text{O} | \phi(x) \phi^\dagger(y) | \text{O} \rangle = \int \frac{dp^3}{(2\pi)^3 2E_p} e^{-ip(x-y)} = \frac{1}{(2\pi)^3} \int \frac{p^2 dp}{2E_p} e^{-ip \cdot (x-y)} \int_0^\pi m\theta d\theta \int_0^\pi d\phi$$

$$= \frac{2\pi}{(2\pi)^3} \int \frac{p^2 dp}{E_p} e^{-ip \cdot (x-y)} \int_0^\pi m\theta d\theta \cdot \frac{d}{ip^2} \frac{d}{d\theta} (e^{-ip^2 \cos\theta})$$

$$= \frac{1}{(2\pi)^2} \int \frac{p^2 dp}{E_p} \frac{d\theta}{ip^2} \frac{d}{d\theta} (e^{-ip^2 \cos\theta})$$

$$x^0 - y^0 = \Delta(x-y) \underset{t \rightarrow \infty}{\sim} e^{-mt}$$

$$\tilde{x} - \tilde{y}$$

(21)

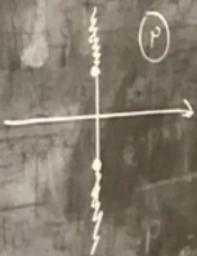
$$|\phi(x)\phi^*(y)|\psi\rangle = \frac{1}{(2\pi)^2} \int \frac{dp}{E_p} \frac{p^2}{i p z} \int_0^\pi d\theta \underbrace{\frac{d}{d\theta} (e^{i p z \cos\theta})}_{\text{...}}$$

$$\begin{aligned} \gamma &:= \Delta(x-y) \underset{t \rightarrow \infty}{\sim} e^{-imt} \quad \left\{ \begin{aligned} &= \frac{1}{(2\pi)^2} \int \frac{dp}{E_p} \frac{p}{z} (e^{i p z} - e^{-i p z}) \\ &= \frac{-i}{(2\pi)^2 z} \int_0^\infty dp \frac{p}{\sqrt{p^2 + m^2}} (e^{i p z} - e^{-i p z}) \quad p = |\vec{p}| \\ &= \frac{-i}{(2\pi)^2 z} \left\{ \int_0^\infty dp \frac{p}{\sqrt{p^2 + m^2}} e^{i p z} - \underbrace{\int_0^\infty dp \frac{p}{\sqrt{p^2 + m^2}} e^{-i p z}}_{\substack{p \rightarrow -p \\ \Rightarrow \int_0^\infty (-dp) \frac{-p}{\sqrt{p^2 + m^2}} e^{i p z} = - \int_0^\infty dp \frac{p}{\sqrt{p^2 + m^2}} e^{i p z}}} \right\} \end{aligned} \right. \end{aligned}$$

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$$\langle 0 | \phi(x) \phi^\dagger(y) | 0 \rangle = \frac{-i}{(2\pi)^2 z} \left\{ \int_0^{+\infty} dp \frac{p}{\sqrt{p^2 + m^2}} e^{-p z} + \int_{-\infty}^0 dp \frac{p}{\sqrt{p^2 + m^2}} e^{ipz} \right\}$$

$$= \frac{-i}{(2\pi)^2 z} \int_{-\infty}^{+\infty} dp \frac{p}{\sqrt{p^2 + m^2}} e^{-ipz} \Rightarrow \text{polarization } \pm \text{ in } m$$



$$S = \frac{-i}{(2\pi)^2 z} \int_m^{\infty} dp \frac{p}{\sqrt{m^2 - p^2}} e^{-ipz} \underset{z \rightarrow \infty}{\sim} e^{-imz} \quad \begin{aligned} x^0 &= y^0 \\ \vec{x} - \vec{y} &= x^1 - y^1 \end{aligned}$$

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